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# Contents

<b>Introduction</b>	<b>3</b>
<b>1 Presentation of the corporation</b>	<b>4</b>
1.1 UL . . . . .	4
1.2 ENSMN . . . . .	4
1.2.1 Logo . . . . .	4
1.3 Lemta . . . . .	4
1.3.1 Logo . . . . .	5
<b>2 Presentation of my work</b>	<b>6</b>
2.1 Teaching activities . . . . .	6
2.2 Research activities . . . . .	7
2.3 Organization . . . . .	7
<b>Conclusion</b>	<b>9</b>
<b>A Curriculum Vitae</b>	<b>10</b>
A.1 Initial training . . . . .	10
A.2 Career . . . . .	10
A.3 Main responsibility . . . . .	10
<b>Bibliography</b>	<b>11</b>

# Introduction

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Please *modify* it !

Nancy - May 7, 2021.

Emmanuel PLAUT.

# Chapter 1

## Presentation of the corporation

I am a Professor at Université de Lorraine (UL) and Researcher in LEMTA. My CV is available in the appendix [A](#).

### 1.1 UL

The *Université de Lorraine* is a great French University that contains ten Engineering Schools. I teach in one of them, the ENSMN.

### 1.2 ENSMN

ENSMN is the *École Nationale Supérieure des Mines de Nancy*.

#### 1.2.1 Logo

I present the logo of ENSMN on figure [1.1](#).



**Fig. 1.1** : Note that the Figure number appears in bold font and that the caption appears with a small font.

### 1.3 Lemta

Lemta is the *Laboratoire d'Énergétique et de Mécanique Théorique et Appliquée*. It depends on the *Centre National de la Recherche Scientifique* (CNRS) and on the UL.



**Fig. 1.2 :** This Figure appears at the top of the page because of the option [t] in the latex file.

### 1.3.1 Logo

I present the logo of Lemta on figure [1.2](#).

# Chapter 2

## Presentation of my work

### 2.1 Teaching activities

Presently one of the most sophisticated topics that I teach concerns *Turbulence modelling*. I give you here a copy of a part of a recent version of the chapter 0 of [Plaut \*et al.\* \(2021\)](#). In this chapter, I recalled the bases of the theory that describes the *dynamics of incompressible newtonian fluids*.

The *eulerian velocity field*  $\mathbf{v}(\mathbf{x}, t)$  is used to describe the flow of *incompressible newtonian fluids*. The *incompressibility* means that the *mass density*  $\rho$  is uniform and constant. Therefore the *mass conservation equation* reads

$$\boxed{\partial_{x_i} v_i = 0}, \quad (2.1)$$

which means that the velocity field  $\mathbf{v}$  is conservative.

The *Cauchy stress tensor*  $\boldsymbol{\sigma}$  determines the surface force  $d^2\mathbf{f}$  exerted on a small surface of area  $d^2A$  and normal unit vector  $\mathbf{n}$  pointing outwards, by the exterior onto the interior, through

$$d^2\mathbf{f} = \boldsymbol{\sigma} \cdot \mathbf{n} d^2A = \sigma_{ij} n_j d^2A \mathbf{e}_i. \quad (2.2)$$

The stresses  $\sigma_{ij}$  are given by the sum of *pressure* and *viscous stresses*,

$$\sigma_{ij} = -p_{\text{static}} \delta_{ij} + 2\eta S_{ij}(\mathbf{v}) \quad (2.3)$$

with  $p_{\text{static}}$  the *static pressure*,  $\eta$  the *dynamic viscosity* of the fluid,  $\mathbf{S}(\mathbf{v})$  the *rate-of-strain tensor* defined as the symmetric part of the velocity gradient, i.e., in components,

$$\boxed{S_{ij}(\mathbf{v}) := \frac{1}{2}(\partial_{x_i} v_j + \partial_{x_j} v_i)}. \quad (2.4)$$

In this equation, the sign  $:=$  means a definition.

The linear momentum equation is the *Navier-Stokes equation*

$$\rho[\partial_t v_i + \partial_{x_j}(v_i v_j)] = \rho g_i + \partial_{x_j} \sigma_{ij} = \rho g_i - \partial_{x_i} p_{\text{static}} + \partial_{x_j}(2\eta S_{ij}(\mathbf{v})), \quad (2.5)$$

where  $g_i$  are the components of the acceleration due to gravity. Usually, a *modified pressure* that includes a gravity term,

$$p = p_{\text{static}} + \rho g Z , \quad (2.6)$$

where  $Z$  is a vertical coordinate, is used, to group the first two terms on the r.h.s. of equation (2.5), which reads therefore

$$\boxed{\rho[\partial_t v_i + \partial_{x_j}(v_i v_j)] = -\partial_{x_i} p + \partial_{x_j}(2\eta S_{ij}(\mathbf{v}))} . \quad (2.7)$$

We will most often use the modified pressure  $p$  instead of the static pressure  $p_{\text{static}}$ .

After dividing by the mass density, we get another form of the *Navier-Stokes equation*,

$$\boxed{\partial_t v_i + v_j \partial_{x_j} v_i = -\frac{1}{\rho} \partial_{x_i} p + \partial_{x_j}(2\nu S_{ij}(\mathbf{v}))} \quad (2.8)$$

with  $\nu = \eta/\rho$  the *kinematic viscosity* of the fluid. We have used the mass conservation equation (2.1),  $\partial_{x_j} v_j = 0$ , to rewrite the nonlinear term on the l.h.s.; this whole l.h.s. is in fact the acceleration of the fluid particle (in the sense of the continuum mechanics) that passes through  $\mathbf{x}$  at time  $t$ ...

## 2.2 Research activities

I try to advance the study of the nonlinear dynamics of extended systems. For instance I developed for this purpose a reformulation of the Reynolds stress tensor  $\boldsymbol{\tau}$  created by pure bidimensional waves at the lowest nonlinear order. This work has been published in [Plaut \*et al.\* \(2008\)](#). It helps to analyze the instability mechanisms in shear flows, by pushing farther the study performed by [Pedlosky \(1987\)](#), or to better understand the form of zonal flows in rotating convection, see e.g. [Morin & Dormy \(2006\)](#). In Cartesian geometry  $xy$ , with  $x$  the periodic direction of the wave, the reformulation reads as follows:

$$\tau_{xx} = -2E_{cx} , \quad \tau_{yy} = -2E_{cy} , \quad \tau_{xy} = \tau_{yy} \tan \alpha . \quad (2.9)$$

There

$$E_{cx} = \frac{1}{2} \langle v_x^2 \rangle_x , \quad E_{cy} = \frac{1}{2} \langle v_y^2 \rangle_x \quad (2.10)$$

are the mean kinetic energies associated to the  $x$  and  $y$  components of the wave velocity field,  $\alpha$  is the angle between the wave separatrices and the  $y$  direction.

Task	Teaching	Research
Theoretical proportion	50%	50%
Experimental observation	$52 \pm 1\%$	$48 \pm 1\%$

**Tab. 2.1 :** Sketch of the partition of my working time between Teaching and Research. Teaching activities also encompass duties linked to my responsibilities, as the Head of an Option of Mines Nancy.

## 2.3 Organization

In table 2.1 I sketch the partition of my working time between Teaching and Research. This partition has evolved with time, I am happy to verify in 2021 an ‘equipartition principle’ !



# Conclusion

I wish you a good play with L<sup>A</sup>T<sub>E</sub>X...

# Appendix A

## Curriculum Vitae

### A.1 Initial training

1989 : École Polytechnique, Palaiseau, France

1992 : DEA (Master) of Theoretical Physics, Paris, France

1996 : PhD in Physics, Université d'Orsay, France

### A.2 Career

1996-1998 : Post-doc at the Theoretical Physics Institute of the University of Bayreuth, Germany

1998-2008 : Assistant professor at École Nationale Supérieure d'Électricité et de Mécanique (ENSEM),  
France

2008-... : Professor at ENSMN, France

1998-... : Researcher in Lemta

### A.3 Main responsibility

Responsibility of the option Energy & Fluid Mechanics at ENSMN.

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